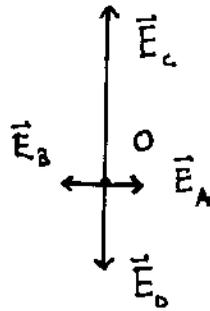


1)

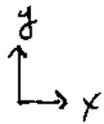
a)

$$-q \quad B \\ (-2d, 0)$$

$$+q \quad D \\ (0, d)$$



$$-q \quad A \\ (2d, 0)$$



$$+2q \quad C \\ (0, -d)$$

$$|\vec{E}_C| > |\vec{E}_D| > |\vec{E}_B| = |\vec{E}_A|$$

Resultant electric field vector:

$$\vec{E}_R = \vec{E}_A + \vec{E}_B + \vec{E}_C + \vec{E}_D$$

b) Let \vec{r}_{A0} be the displacement from A to O.

Since $|\vec{r}_{A0}| = |\vec{r}_{B0}|$ and $q_A = q_B$, we have that the electric field contributions from A and B are equal and opposite and thus cancel.

It remains to calculate the contributions from C and D:

$$\vec{E}_C = \frac{k(2q)}{d^2} \hat{r}_{C0} = \frac{2kq}{d^2} \hat{y} \quad \text{where } k = \frac{1}{4\pi\epsilon_0}$$

$$\vec{E}_D = \frac{k(q)}{d^2} \hat{r}_{D0} = \frac{kq}{d^2} (-\hat{y})$$

$$\Rightarrow \vec{E}_{\text{TOT}} = \frac{2kq}{d^2} \hat{y} - \frac{kq}{d^2} \hat{y} = \boxed{\frac{kq}{d^2} \hat{y}}$$

c) In principle, one could obtain the electric field from the electric potential $V(\vec{r})$. Recall that the electric field is just the slope of the electric potential function. In one dimension,

$$E(x) = -\frac{dV(x)}{dx}$$

Using the superposition principle, $V(\vec{r})$ can be calculated as $V(\vec{r}) = V_A(\vec{r}) + V_B(\vec{r}) + V_C(\vec{r}) + V_D(\vec{r})$

$$\text{where } V_A(\vec{r}) = \frac{kq_A}{|\vec{r}_A|}$$

is the potential from a point charge.

2: a) There are two plates, each of which generates an electric field which can be calculated with Gauss' Law. The total \vec{E} -field in each region follows from superposition principle.



$$\begin{aligned}\vec{E}_{in} &= \vec{E}_+ + \vec{E}_- = \frac{\sigma}{2\epsilon_0} (-\hat{j}) + \frac{\sigma}{2\epsilon_0} (-\hat{j}) = \frac{\sigma}{\epsilon_0} (-\hat{j}) \\ &= \frac{Q}{A\epsilon_0} (-\hat{j});\end{aligned}$$

$$\vec{E}_{out} = \vec{E}_+ + \vec{E}_- = 0.$$

b) $\Delta V = -\int \vec{E} \cdot d\vec{r}$. For uniform field, the magnitude is $\Delta V = E \cdot d = \frac{Qd}{A\epsilon_0}$.

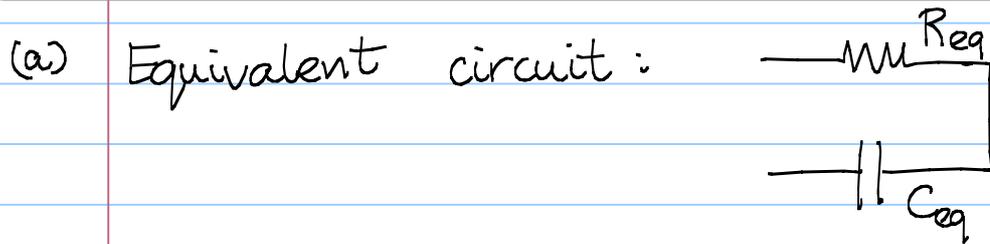
c) The definition of capacitance gives

$$C = \frac{Q}{\Delta V} = \frac{A\epsilon_0}{d}.$$

8b Fall '09 Bordel Midterm 1 Problem 3 Solution

Note Title

10/8/2009



$$R_{eq} = \frac{R_1 R_2}{R_1 + R_2} + R_3 = \frac{R(3R)}{R + 3R} + \frac{R}{2} \Rightarrow \boxed{R_{eq} = \frac{5}{4} R}$$

$$C_{eq} = C_1 + C_2 \Rightarrow \boxed{C_{eq} = 3C}$$

(b) $V_{Req} = \mathcal{E} - V_{Ceq}$ (by Kirchoff's voltage-loop law)

$$\Rightarrow V_{Req}^{max} = \mathcal{E} \Rightarrow \boxed{P_{Req}^{max} = \frac{(V_{Req}^{max})^2}{R_{eq}} = \frac{4\mathcal{E}^2}{5R}}$$

(c) $V_{Ceq}^{max} = \mathcal{E} \Rightarrow \boxed{U_{Ceq}^{max} = \frac{1}{2} C_{eq} (V_{Ceq}^{max})^2 = \frac{3}{2} C\mathcal{E}^2}$

(d)(i) $I = C \frac{dV}{dt} \Rightarrow V_{Ceq}$ must be continuous \Rightarrow since $V_{Ceq} = 0$ right before the battery is connected, then it is 0 right after \Rightarrow all the voltage \mathcal{E} is initially dropped across $R_{eq} \Rightarrow \boxed{I = \frac{\mathcal{E}}{R_{eq}} = \frac{4\mathcal{E}}{5R}$ (short time)

Note: this is the current through the resistor, but since it is in series with the capacitor, this is also the current through the capacitor.

Since $V_{Ceq} = 0$ initially, $\boxed{Q_{Ceq} = 0}$ (short time)

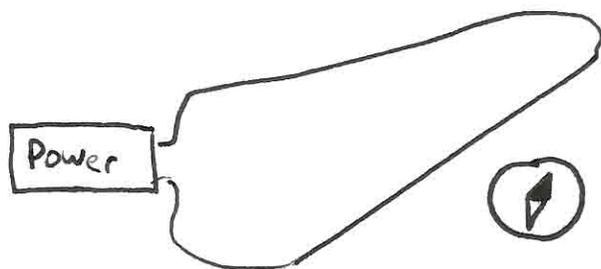
(ii) As $t \rightarrow \infty$, the circuit settles down so that $\frac{dV}{dt} = 0$
 $\Rightarrow \boxed{I = 0}$ (long time) $\Rightarrow V_{Req} = IR_{eq} = 0 \Rightarrow V_{Ceq} = \mathcal{E}$

$$\Rightarrow \boxed{Q_{Ceq} = C_{eq} \mathcal{E} = 3C\mathcal{E}} \text{ (long time)}$$

Problem 4 (10 Points)

a) Oersted's experiment used a compass to demonstrate the existence of a magnetic field created by a current.

To recreate this experiment, you would want a compass, some wire, and some current source: a battery, power supply, etc.



The basic idea is to show that the compass (itself a small magnet) is affected by the presence of a current-carrying wire. Thus, you could turn the power supply off, move the ~~the~~ compass around the wire, and see no effect. Then, after turning the power supply on, the compass will be suddenly affected by the presence of the wire, indicating that a magnetic field is created. Other possible answers include using a ~~small~~ magnet in place of the ~~the~~ compass, iron filings, ~~or~~ using a solenoid...

b)

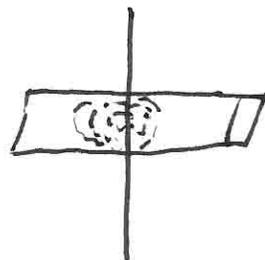
To visualize the field lines, you could:

- use a compass ~~to~~ to trace out the direction of \vec{B} as you move it around the wire

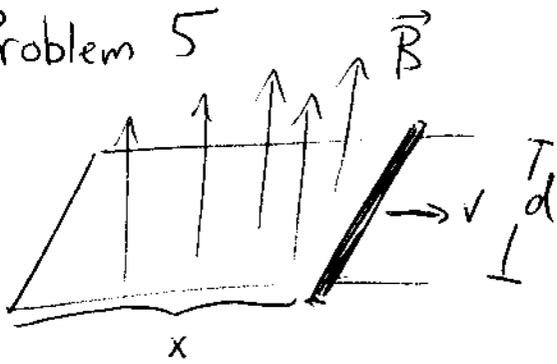
- use iron filings to visualize the field:
 - viewing paper

There are ~~many~~ other correct answers as well.

(Right hand rule, etc)

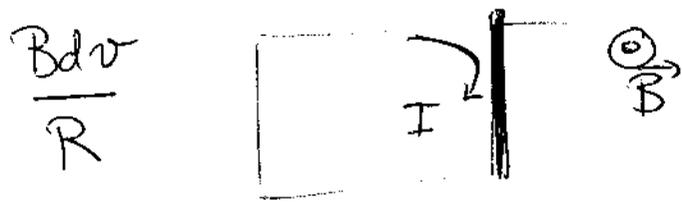


Problem 5



(a) $\Phi_B = \vec{B} \cdot \vec{A} = BA = Bdx$, $x(t) = vt$, $\Phi_B(t) = Bdv t$

(b) $\mathcal{E}_{emf} = -\frac{d\Phi_B}{dt} = -Bdv$. By Lenz's law, the induced current generates a field anti-parallel to \vec{B} . By RHR, current flows clockwise as viewed from above. It has magnitude



(c) $\vec{F} = I\vec{L} \times \vec{B}$. $|\vec{F}| = \frac{Bdv}{R} \cdot dB = \frac{(Bd)^2 v}{R}$

By RHR, \vec{F} directed opposite to \vec{v} . This agrees with Lenz's law since the force opposes the motion and thereby resists the increase in flux lines through the frame.

(d) $W = |\vec{F}|L = \frac{(Bd)^2 v}{R} L$ Alternatively,

$P = I^2 R = \left(\frac{Bdv}{R}\right)^2 R = \frac{(Bd)^2 v^2}{R} = \frac{\Delta E}{\Delta t}$, $\Delta E = \frac{(Bd)^2 v}{R} (v\Delta t)$

$= \frac{(Bd)^2 v}{R} L$