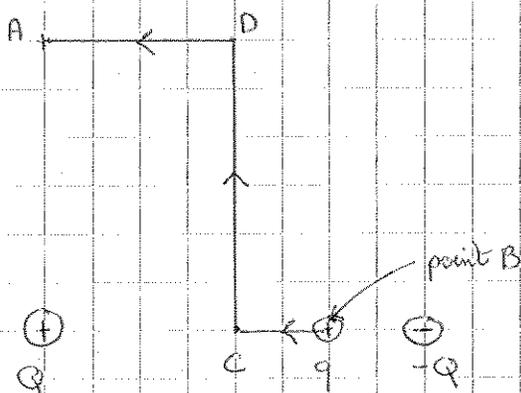


Physics 8B - Fall 2011 - Mid Term 1, October 5, 2011

Problem 1



Key Idea: Static electric fields are conservative.

Thus the work done on the test charge is independent of the path!

Be W the work done on the test charge, U_A, U_B the potential energy of the test charge at points A, B.

$$W = + \Delta U = U_A - U_B = q(V_A - V_B) \text{ where } V \text{ is the electric potential}$$

By the principle of superposition, the electric potential at one point is the sum of the potentials created by each charge:

$$V_B = V_B^+ + V_B^- \text{ where } V_B^+ \text{ is the electric potential created by the charge } +Q \text{ at B.}$$

Thus: $W = q(V_A^+ + V_A^- - V_B^+ - V_B^-)$

Now, points A and B are at the same distance of the charge $+Q$

Hence $V_B^+ = V_A^+$

and $W = q(V_A^- - V_B^-)$

→ The work does not depend on the value of the positive charge.

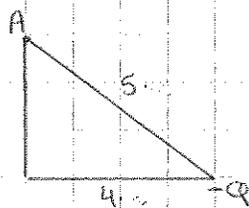
We can actually consider that the positive charge is not there!

The electric potential created by a point charge $-Q$ is

$$V = \frac{-Q}{4\pi\epsilon_0 r} = \frac{-RQ}{r}$$

Thus $V_B^- = -\frac{RQ}{1} = -RQ$

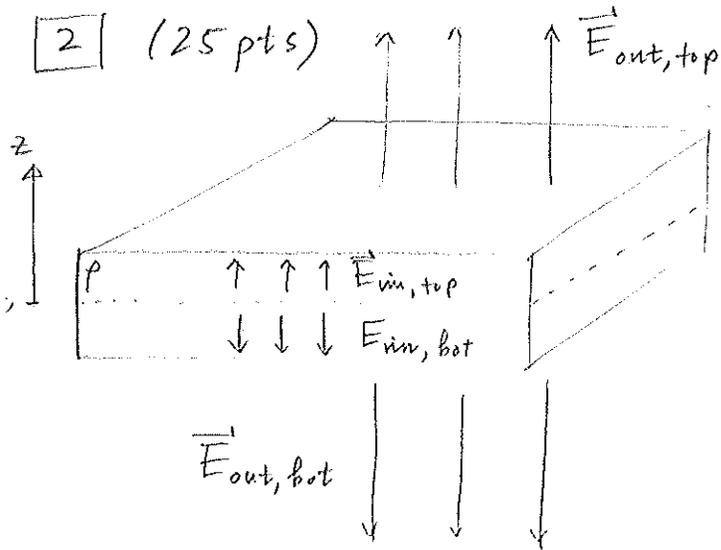
and $V_A^- = -\frac{RQ}{5}$ since: 3



The work done is:

$$W = \frac{4}{5} RqQ \quad \left(R = \frac{1}{4\pi\epsilon_0} \right)$$

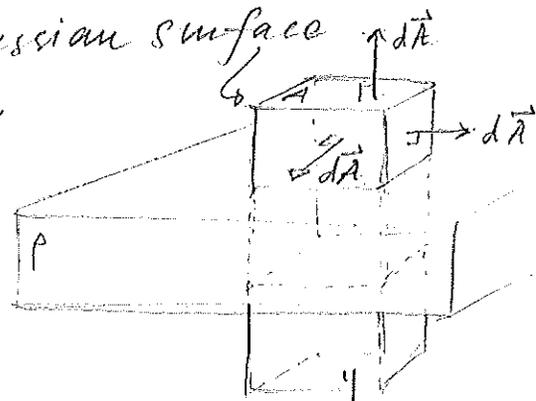
2 (25 pts)



By symmetry, we expect \vec{E} to point in the \hat{z} direction for $z > 0$ and in the $-\hat{z}$ direction for $z < 0$. We also expect $\vec{E}(z=0) = \vec{0}$, as well as $\vec{E}_{in,top} = -\vec{E}_{in,bot}$ and $\vec{E}_{out,top} = -\vec{E}_{out,bot}$.

Thus, we only need to find $|\vec{E}_{out}|$ and $|\vec{E}_{in}|$.

To find $|\vec{E}_{out}|$, consider the following Gaussian surface. Begin by writing Gauss' Law and evaluating the left-hand side and right-hand side separately.



$$\oint \vec{E} \cdot d\vec{A} = \frac{Q_{enclosed}}{\epsilon_0}$$

integral over closed Gaussian surface

$$\begin{aligned} \oint \vec{E} \cdot d\vec{A} &= \int_{top} \vec{E} \cdot d\vec{A} + \int_{bottom} \vec{E} \cdot d\vec{A} + \int_{left} \vec{E} \cdot d\vec{A} + \int_{right} \vec{E} \cdot d\vec{A} + \int_{front} \vec{E} \cdot d\vec{A} + \int_{back} \vec{E} \cdot d\vec{A} \\ &= |\vec{E}_{out}| \left(\int_{top} dA \underbrace{\cos 0}_{=1} + \int_{bottom} dA \underbrace{\cos 0}_{=1} + \int_{left} dA \underbrace{\cos \frac{\pi}{2}}_{=0} + \int_{right} dA \underbrace{\cos \frac{\pi}{2}}_{=0} + \int_{front} dA \underbrace{\cos \frac{\pi}{2}}_{=0} + \int_{back} dA \underbrace{\cos \frac{\pi}{2}}_{=0} \right) \\ &= |\vec{E}_{out}| (A + A) = 2A |\vec{E}_{out}| \end{aligned}$$

recall that $\rho = \frac{dQ}{dV}$:

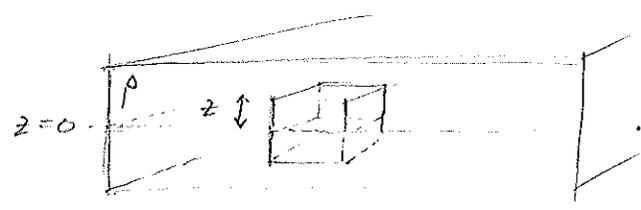
$$\frac{Q_{encl}}{\epsilon_0} = \frac{1}{\epsilon_0} \int dQ = \frac{1}{\epsilon_0} \int \rho dV = \frac{\rho}{\epsilon_0} \int dV = \frac{\rho}{\epsilon_0} Ad$$

Thus, Gauss' Law gives: $2A |\vec{E}_{out}| = \frac{\rho}{\epsilon_0} Ad \Rightarrow \begin{cases} \vec{E}_{out,top} = \frac{\rho d}{2\epsilon_0} \hat{z} \\ \vec{E}_{out,bot} = -\frac{\rho d}{2\epsilon_0} \hat{z} \end{cases}$

Note that outside the slab, electric field is independent of z , as we have seen in an infinite sheet of charge.

To find $|\vec{E}_{in}|$, consider another Gaussian surface:

$$\oint \vec{E} \cdot d\vec{A} = \int_{top} \vec{E} \cdot d\vec{A} + \int_{bottom} \vec{E} \cdot d\vec{A} + \int_{left} \vec{E} \cdot d\vec{A} + \int_{right} \vec{E} \cdot d\vec{A} + \int_{front} \vec{E} \cdot d\vec{A} + \int_{back} \vec{E} \cdot d\vec{A}$$



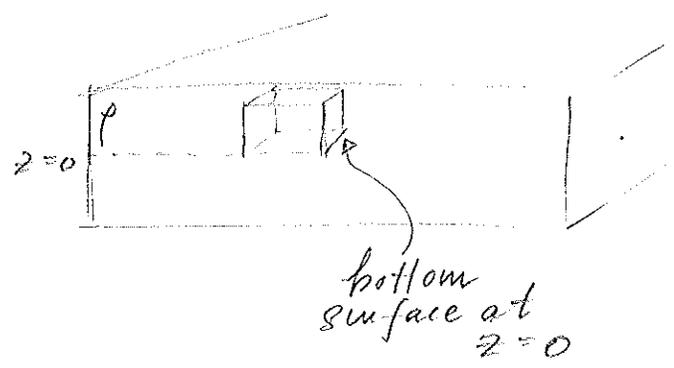
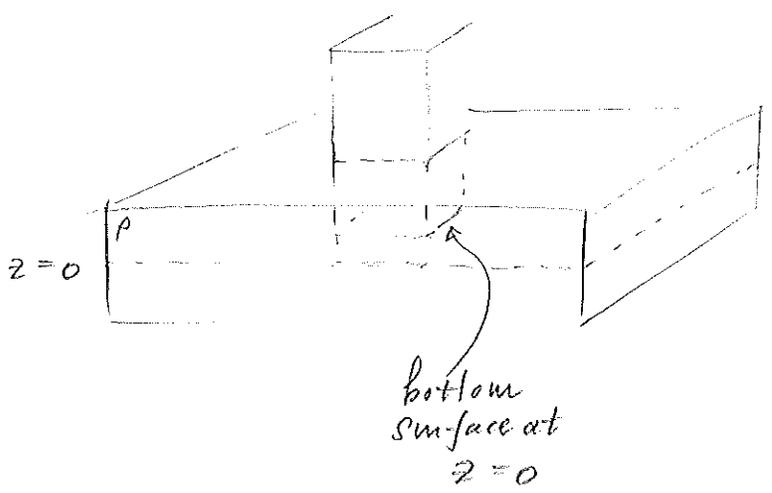
$$= |\vec{E}_{in}| \left(\int_{top} dA \cos 0 + \int_{bottom} dA \cos 0 + \int_{left} dA \cos \frac{\pi}{2} + \int_{right} dA \cos \frac{\pi}{2} + \int_{front} dA \cos \frac{\pi}{2} + \int_{back} dA \cos \frac{\pi}{2} \right)$$

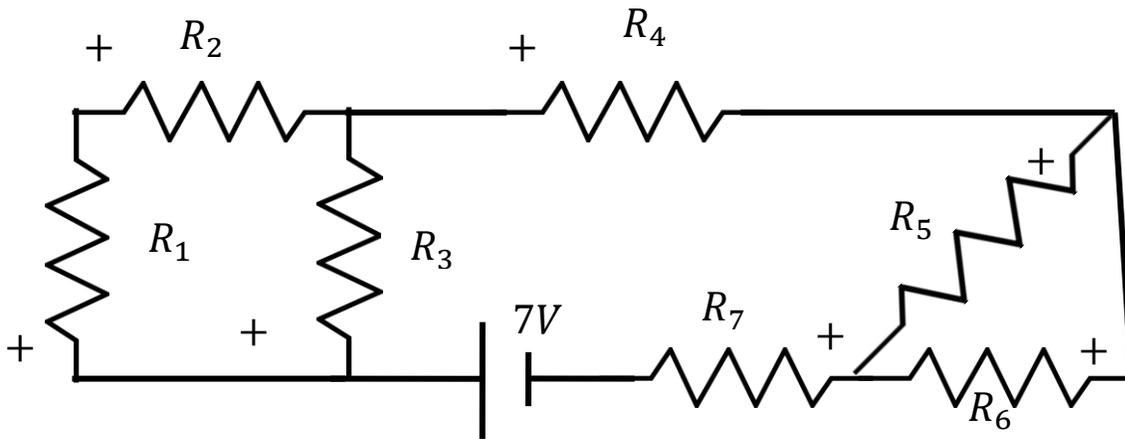
$$= |\vec{E}_{in}| (A + A) = 2A |\vec{E}_{in}|$$

$$\frac{Q_{encl}}{\epsilon_0} = \frac{1}{\epsilon_0} \int dQ = \frac{\rho}{\epsilon_0} \int dV = \frac{\rho}{\epsilon_0} 2zA$$

$$\therefore \text{Gauss' Law yields } 2A |\vec{E}_{in}| = \frac{\rho}{\epsilon_0} 2zA \Rightarrow \begin{cases} \vec{E}_{in, top} = \frac{\rho z}{\epsilon_0} \hat{z} \\ \vec{E}_{in, bottom} = -\frac{\rho z}{\epsilon_0} \hat{z} \end{cases}$$

Note that we could've also chosen the following Gaussian surfaces and gotten the same answers (try it!):





Higher potential is denoted by + on one end of each of the resistors

$$R_{12} = R_1 + R_2 = 12\Omega$$

$$R_{123} = \frac{1}{\frac{1}{R_{12}} + \frac{1}{R_3}} = \frac{1}{\frac{1}{12} + \frac{1}{6}} \Omega = 4\Omega$$

$$R_{56} = \frac{1}{\frac{1}{R_5} + \frac{1}{R_6}} = 5\Omega$$

$$R_{eq} = R_{123} + R_4 + R_5 + R_{56} + R_7 = 14\Omega$$

$$V = I_{total} R_{eq} \rightarrow 7V = I_{total} \times 14\Omega \rightarrow I_{total} = \frac{1}{2} A$$

$$V_4 = I_{total} \times R_4 = \frac{1}{2} V$$

$$V_7 = I_{total} \times R_7 = 2V$$

$$V_5 = V_6 = \frac{I_{total}}{2} \times R_5 = \frac{5}{2} V$$

$$V_3 = V_1 + V_2 = 7 - V_4 - V_5 - V_7 = 2V$$

$$I_1 = I_2 \text{ and } R_1 = R_2 \rightarrow V_1 = V_2 = \frac{1}{2} V_3 = 1V$$

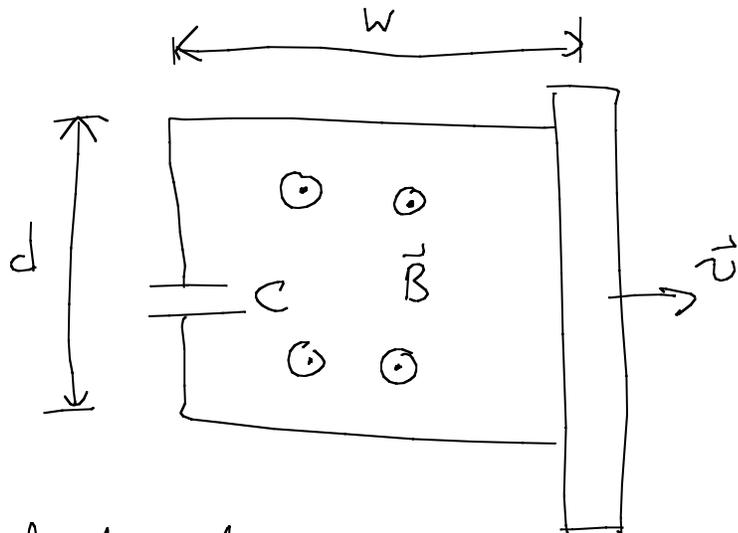


$$V_1 = V_2 = 1V, V_3 = 2V, V_4 = \frac{1}{2} V, V_5 = V_6 = \frac{5}{2} V, V_7 = 2V$$

Problem 4 solution (Phys 8B Section I exam)

Note Title

10/15/2011



Induced EMF:

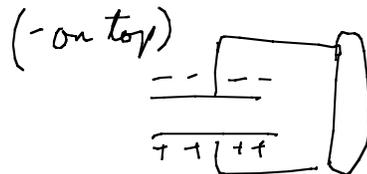
$$|\mathcal{E}| = \left| \frac{d\Phi_B}{dt} \right| \text{ so we need to solve for flux and differentiate it.}$$

$$\left. \begin{array}{l} w = w_0 + vt \\ d = \text{const} \end{array} \right\} A = wd \left. \begin{array}{l} B = \text{const} \end{array} \right\} \begin{aligned} \Phi_B &= BA \\ &= Bwd \\ &= Bd(w_0 + vt) \end{aligned}$$

$$\left| \frac{d\Phi_B}{dt} \right| = Bd v = |\mathcal{E}|$$

For a capacitor, $Q = CV$, and since there is no resistance, it is charged instantly

$$Q = C B v d$$



To find the sign of charge on each plate, either use Lenz's law, or $\vec{F} = q(\vec{v} \times \vec{B})$

⊕ charge: \vec{B} into page, \vec{v} to the right, \vec{F} down

⊖ charge: \vec{B} into page, \vec{v} to the right, \vec{F} up